## Fall Term 2002 Introduction to Plasma Physics I

## 22.611J, 6.651J, 8.613J

Problems for study

1. Collision Operator Properties. The Coulomb collision operator for like-particle collisions in Landau form is given by,

$$\mathcal{C}(f,f) = \Gamma \frac{\partial}{\partial \mathbf{v}} \cdot \int d^3 v' \mathbf{U}(\mathbf{v} - \mathbf{v}') \cdot \left(\frac{\partial}{\partial \mathbf{v}} - \frac{\partial}{\partial \mathbf{v}'}\right) f(\mathbf{v}) f(\mathbf{v}')$$

with,

$$\mathbf{U}(\mathbf{v} - \mathbf{v}') = \frac{1}{|\mathbf{v} - \mathbf{v}'|} \left( \mathbf{I} - \frac{(\mathbf{v} - \mathbf{v}') (\mathbf{v} - \mathbf{v}')}{|\mathbf{v} - \mathbf{v}'|^2} \right)$$
$$\Gamma = \frac{2\pi q^4 \ln \Lambda}{m^2}$$

Prove the conservation laws for the collision operator:

Particle Conservation
$$0 = \int d^3 v C(f, f)$$
Momentum Conservation $0 = \int d^3 v m \mathbf{v} C(f, f)$ Energy Conservation $0 = \int d^3 v \frac{1}{2} m v^2 C(f, f)$ 

2. *Two Species Collisions*: Show that the conservation laws also hold for collision between two species, but that the sum over species is now required:

Particle Conservation
$$0 = \int d^3 v \mathcal{C}_{\alpha\beta}(f_\alpha, f_\beta)$$
Momentum Conservation $0 = \int d^3 v \left[m_\alpha \mathbf{v} \mathcal{C}_{\alpha\beta}(f_\alpha, f_\beta) + m_\beta \mathbf{v} \mathcal{C}_{\beta\alpha}(f_\beta, f_\alpha)\right]$ Energy Conservation $0 = \int d^3 v \left[\frac{1}{2}m_\alpha v^2 \mathcal{C}_{\alpha\beta}(f_\alpha, f_\beta) + \frac{1}{2}m_\beta v^2 \mathcal{C}_{\beta\alpha}(f_\beta, f_\alpha)\right]$  $0 = \int d^3 v \left[\frac{1}{2}m_\alpha v^2 \mathcal{C}_{\alpha\beta}(f_\alpha, f_\beta) + \frac{1}{2}m_\beta v^2 \mathcal{C}_{\beta\alpha}(f_\beta, f_\alpha)\right]$ 

Note that particle's are conserved by *each* collision operator, while momentum and energy is only conserved as a species sum. Make a statement of the physical meaning of the two separate conservation laws for each of momentum and energy. Recall that the collision operator between species,  $\alpha$ , and,  $\beta$ , is given by,

$$\mathcal{C}_{\alpha\beta}(f_{\alpha}, f_{\beta}) = \Gamma_{\alpha\beta} \frac{\partial}{\partial \mathbf{v}} \cdot \int d^{3}v' \mathbf{U}(\mathbf{v} - \mathbf{v}') \cdot \left(\frac{m_{\beta}}{m_{\alpha}} \frac{\partial}{\partial \mathbf{v}} - \frac{\partial}{\partial \mathbf{v}'}\right) f_{\alpha}(\mathbf{v}) f_{\beta}(\mathbf{v}')$$

$$\Gamma_{\alpha\beta} = \frac{2\pi Z_{\alpha}^{2} Z_{\beta}^{2} e^{4} \ln \Lambda}{m_{\alpha} m_{\beta}}$$

and the tensor, U, defined in terms of particle velocities, has the same definition as above.

3. *Temperature Equilibration*: Show that temperature equilibration proceeds according to a term of the form,

$$\frac{\partial}{\partial t}\frac{3}{2}nT_e = -\nu_{ei}\frac{m_e}{m_i}\sqrt{\frac{2}{\pi}n\left(T_e - T_i\right)}$$

by taking the energy moment,

$$\int d^3v \frac{1}{2} m_e v^2 \mathcal{C}_{ei}\left(f_e^{\max}, f_i^{\max}\right)$$

You may use the expanded form of,  $C_{ei}$ , obtained by assuming a Maxwellian distribution for the ions,

$$\mathcal{C}_{ei} \simeq \mathcal{C}_{ei}^{L}\left(f_{e}\right) + \mathcal{C}_{ei}^{E}\left(f_{e}\right)$$

with the Lorentz collision operator,

$$\mathcal{C}_{ei}^{L}(f_{e}) \equiv \nu_{ei} \frac{v_{e}^{3}}{v^{3}} \mathcal{L}(f_{e})$$
$$\mathcal{L} \equiv \frac{1}{2} \frac{\partial}{\partial \mu} (1 - \mu^{2}) \frac{\partial}{\partial \mu}$$

and the energy exchange operator,

$$\mathcal{C}_{ei}^{E}\left(f_{e}\right) = \nu_{ei}\frac{m_{e}}{2m_{i}}v_{e}^{3}\left[\frac{1}{v^{2}}\frac{\partial}{\partial v} + \frac{T_{i}}{m_{e}v^{2}}\frac{\partial}{\partial v}\frac{1}{v}\frac{\partial}{\partial v}\right]f_{e}$$

The electron-ion collision frequency defined as,

$$\nu_{ei} \equiv \frac{4\pi n e^4 \ln \Lambda}{m_e v_e^3}$$

and we have used the convention,  $v_e = \sqrt{T_e/m_e}$ , for the thermal velocity. n.b. different from convention in Landau problem to make expressions cleaner. . .

What is the relative rate of angle scattering vs. thermalization for electrons scattering off ions?