## Recitation 2: Stress/Strain Transformations and Mohr's Circle

### 2.1 General Transformation Rules

### 2.1.1 2D Vector

Consider a vector $\boldsymbol{v}$ in the $\left(x_{1}, x_{2}\right)$ coordinate system. A new coordinate system $\left(x_{1}^{\prime}, x_{2}^{\prime}\right)$ is obtained by rotating the old coordinate system by angle $\theta$. Find the components of $\boldsymbol{v}$ in the new coordinate system.


From geometry, we have:

$$
\begin{align*}
x_{1}^{\prime} & =x_{1} \cos \theta+x_{2} \sin \theta \\
& =x_{1} \cos \theta+x_{2} \cos \left(\frac{\pi}{2}-\theta\right)  \tag{2.1}\\
& =x_{1} L_{11}+x_{2} L_{12} \\
x_{2}^{\prime} & =-x_{1} \sin \theta+x_{2} \cos \theta \\
& =x_{1} \cos \left(\frac{\pi}{2}+\theta\right)+x_{2} \cos \theta  \tag{2.2}\\
& =x_{1} L_{21}+x_{2} L_{22}
\end{align*}
$$

where $L_{i j}$ is the direction cosine of the basis vectors:

$$
\begin{equation*}
L_{11}=\cos \left(\boldsymbol{e}_{1}^{\prime}, \boldsymbol{e}_{1}\right), L_{12}=\cos \left(\boldsymbol{e}_{1}^{\prime}, \boldsymbol{e}_{2}\right), L_{21}=\cos \left(\boldsymbol{e}_{2}^{\prime}, \boldsymbol{e}_{1}\right), L_{22}=\cos \left(\boldsymbol{e}_{2}^{\prime}, \boldsymbol{e}_{2}\right) \tag{2.3}
\end{equation*}
$$

### 2.1.2 General Vector Transformation

A vector is a physical quantity, independent of the coordinate system. But its scalar components DO depend on the coordinate system.

$$
\begin{equation*}
\boldsymbol{v}=v_{i} \boldsymbol{e}_{i}=v_{i}^{\prime} \boldsymbol{e}_{i}^{\prime} \tag{2.4}
\end{equation*}
$$

The scalar components $v_{i}^{\prime}$ in the new coordinate system can be expressed in terms of the scalar components $v_{i}$ in the original coordinate system. Multiplying the above equation by $\boldsymbol{e}_{j}^{\prime}$ gives:

$$
\begin{align*}
v_{i} \boldsymbol{e}_{i} \cdot \boldsymbol{e}_{j}^{\prime} & =v_{i}^{\prime} \boldsymbol{e}_{i}^{\prime} \cdot \boldsymbol{e}_{j}^{\prime} \\
& =v_{i}^{\prime} \delta_{i^{\prime} j^{\prime}}  \tag{2.5}\\
& =v_{j}^{\prime}
\end{align*}
$$

The indices $i, j$ are arbitrary and may be reversed to give:

$$
\begin{equation*}
v_{i}^{\prime}=v_{j} \boldsymbol{e}_{i}^{\prime} \cdot \boldsymbol{e}_{j} \tag{2.6}
\end{equation*}
$$

Introduce the direction cosine tensor

$$
\begin{equation*}
L_{i j} \equiv \boldsymbol{e}_{i}^{\prime} \cdot \boldsymbol{e}_{j}=\cos \left(\boldsymbol{e}_{i}^{\prime}, \boldsymbol{e}_{j}\right) \quad(i, j=1,2,3) \tag{2.7}
\end{equation*}
$$

then we can write $v_{i}^{\prime}$ in index form

$$
\begin{equation*}
v_{i}^{\prime}=L_{i j} v_{j} \tag{2.8}
\end{equation*}
$$

and in matrix form

$$
\begin{equation*}
\boldsymbol{v}^{\prime}=[L] \boldsymbol{v} \tag{2.9}
\end{equation*}
$$

where

$$
\boldsymbol{v}^{\prime}=\left[\begin{array}{c}
v_{1}^{\prime}  \tag{2.10}\\
v_{2}^{\prime} \\
v_{3}^{\prime}
\end{array}\right], \text { and } \boldsymbol{v}=\left[\begin{array}{l}
v_{1} \\
v_{2} \\
v_{3}
\end{array}\right]
$$

Similarly, the inverse transform

$$
\begin{align*}
& v_{i}=L_{j i} v_{j}^{\prime}  \tag{2.11}\\
& \boldsymbol{v}=[L]^{T} \boldsymbol{v}^{\prime} \tag{2.12}
\end{align*}
$$

### 2.1.3 Tensor Transformations

Using the same definition for $L_{i j}$ defined above, the components of a tensor can be transformed to a new coordinate system by:

$$
\begin{equation*}
A_{i j}^{\prime}=L_{i k} L_{j l} A_{k l} \tag{2.13}
\end{equation*}
$$

or in matrix notation:

$$
\begin{equation*}
\left[\boldsymbol{A}^{\prime}\right]=[\boldsymbol{L}][\boldsymbol{A}][\boldsymbol{L}]^{T} \tag{2.14}
\end{equation*}
$$

The inverse transformation can be found by:

$$
\begin{gather*}
A_{k l}=L_{i k} L_{j l} A_{i j}^{\prime}  \tag{2.15}\\
{[\boldsymbol{A}]=[\boldsymbol{L}]^{T}\left[\boldsymbol{A}^{\prime}\right][\boldsymbol{L}]} \tag{2.16}
\end{gather*}
$$

Note that the transformation matrix $[\boldsymbol{L}]$ is orthogonal, meaning its transpose is equal to its inverse.

$$
\begin{equation*}
[\boldsymbol{L}][\boldsymbol{L}]^{T}=\boldsymbol{I} \tag{2.17}
\end{equation*}
$$

### 2.2 Eigenvalues and Eigenvecors

If there exists a vector $\boldsymbol{n}_{i}$ and a scalar $\lambda_{i}$ for an arbitrary tensor $[\boldsymbol{A}]$ such that

$$
\begin{equation*}
[\boldsymbol{A}] \boldsymbol{n}_{i}=\lambda_{i} \boldsymbol{n}_{i} \tag{2.18}
\end{equation*}
$$

then $\lambda_{i}$ and $\boldsymbol{n}_{i}$ are eigenvalues and eigenvectors of tensor $[\boldsymbol{A}]$, respectively. The three eigenvectors are mutually perpendicular, and form a new coordinate system.

The "eigenvalue problem" for tensor $[\boldsymbol{A}]$ can be written as:

$$
\begin{equation*}
\left(\boldsymbol{A}-\lambda_{i} \boldsymbol{I}\right) \boldsymbol{n}_{i}=\mathbf{0} \tag{2.19}
\end{equation*}
$$

This equation has solutions $\boldsymbol{n}_{i} \neq \mathbf{0}$ only if

$$
\begin{equation*}
\operatorname{det}\left(\boldsymbol{A}-\lambda_{i} \boldsymbol{I}\right)=0 \tag{2.20}
\end{equation*}
$$

This equation can be solved for $\lambda_{i}$, then $\lambda_{i}$ substituted into the eigenvalue problem to solve for $\boldsymbol{n}_{i}$.

For the stress (strain) tensor, the eigenvalues represent principal stresses (strains), and eigenvectors represent principal axes (i.e., faces with zero shear stress (strain)).

## Example

Find the principal stresses and principal axes for $[\sigma]=\left[\begin{array}{cc}80 & 30 \\ 30 & 40\end{array}\right]$.
Solution: The given stress tensor can be represented graphically by


The principal stresses are the eigenvalues $\left(\lambda_{i}\right)$ of the stress tensor, and are found by solving:

$$
\begin{align*}
& \operatorname{det}(\sigma-\lambda \boldsymbol{I})=0  \tag{2.21a}\\
& \operatorname{det}\left[\begin{array}{cc}
80-\lambda & 30 \\
30 & 40-\lambda
\end{array}\right]=0  \tag{2.21b}\\
&(80-\lambda)(40-\lambda)-30^{2}=0  \tag{2.21c}\\
& \lambda^{2}-120 \lambda+2300=0  \tag{2.21d}\\
& \lambda_{1}=96.05 \text { and } \lambda_{2}=23.95 \tag{2.21e}
\end{align*}
$$

The principal directions are the eigenvectors, found by substituting the eigenvalues into the original eigenvalue problem:

$$
\begin{equation*}
\left(\sigma-\lambda_{i} \boldsymbol{I}\right) \boldsymbol{n}_{i}=\mathbf{0} \tag{2.22}
\end{equation*}
$$

For $\lambda_{1}=96.05 \mathrm{MPa}$,

$$
\begin{gather*}
{\left[\begin{array}{cc}
80-96.05 & 30 \\
30 & 40-96.05
\end{array}\right]\left\{\begin{array}{l}
x_{1} \\
y_{1}
\end{array}\right\}=\left\{\begin{array}{l}
0 \\
0
\end{array}\right\}}  \tag{2.23}\\
(80-96.05) x_{1}+30 y_{1}=0 \tag{2.24}
\end{gather*}
$$

Any $\left(x_{1}, y_{1}\right)$ that satisfies this equation is an eigenvector. Therefore, we can choose $y_{1}=1$ and solve for $x_{1}$, then normalize the vector as follows:

$$
\begin{align*}
& (80-96.05) x_{1}+30=0  \tag{2.25}\\
& x_{1}=\frac{30}{16.05}  \tag{2.26}\\
& \boldsymbol{n}_{1}=\frac{1}{\sqrt{\left(\frac{30}{16.05}\right)^{2}+1^{2}}}\left\{\begin{array}{c}
\frac{30}{16.05} \\
1
\end{array}\right\}=\left\{\begin{array}{l}
0.88 \\
0.47
\end{array}\right\} \tag{2.27}
\end{align*}
$$

Similarly, the second eigenvector is found to be:

$$
\boldsymbol{n}_{2}=\left\{\begin{array}{c}
-0.47  \tag{2.28}\\
0.88
\end{array}\right\}
$$

As a check, it is clear that the two eigenvectors are perpendicular $\left(\boldsymbol{n}_{1} \cdot \boldsymbol{n}_{2}=0\right)$.

### 2.3 Mohr's Circle

Mohr's circle (named after Otto Mohr (1835-1918)) is a graphical technique to transform stress (strain) from one coordinate system to another, and to find maximum normal and shear stresses (strains).

## Constructing 2D Mohr's Circle:

1. Establish a rectangular coordinate system with $x=$ normal stress, $y=$ shear stress. Scales must be identical.
2. Plot stresses for 2 orthogonal adjacent faces (values from the original stress (strain) tensor).
3. Connect the 2 points to find center of the circle, C .
4. Draw circle through 2 points with center C.
5. Principal stresses (strains) are values where the circle crosses the x-axis.
6. Max shear stress (strain) is max $y$-value on the circle.

Sign convention for Mohrs circle: Positive shear stress on a face causes clockwise rotation of the unit square.

The stress state for a face rotated an angle $\theta$ from an original coordinate axis may be found by rotating an angle $2 \theta$ on Mohrs circle in the same direction from the original coordinate axis.

## Example ${ }^{1}$

Solve the above example using Mohr's circle.
Solution: The given stress tensor is represented graphically by

[^0]

The stress state on the positive $x$-face is $\sigma=+80 \mathrm{MPa}$, and $\tau=-30 \mathrm{MPa}$ (because of the sign convention defined above).

The stress state on the positive $y$-face is $\sigma=+40 \mathrm{MPa}$, and $\tau=+30 \mathrm{MPa}$.

The Mohrs circle for the given stress state is as shown:


The center, C, is located at $(40+80) / 2=60 \mathrm{MPa}$ on the $\sigma$ axis.

The principal stresses are represented by $A_{1}$ and $B_{1}$. The coordinates of those points are
found from geometry as:

$$
\begin{equation*}
\sigma_{1,2}=60 \pm \sqrt{(80-60)^{2}+30^{2}} \mathrm{MPa} \tag{2.29}
\end{equation*}
$$

or

$$
\begin{equation*}
\sigma_{1}=96.05 \mathrm{MPa} \text { and } \sigma_{2}=23.95 \mathrm{MPa} \tag{2.30}
\end{equation*}
$$

The principal directions are found by:

$$
\begin{gather*}
2 \theta_{\mathrm{p}}=\arctan \frac{30}{(80-60)}=56.3^{\circ}  \tag{2.31a}\\
\theta_{\mathrm{p}}=28.15^{\circ} \tag{2.31b}
\end{gather*}
$$

The principal stress state is as shown below:


## 3D Mohrs Circle

To draw Mohrs circle for a general 3D stress state, the principal stresses and directions must first be evaluated (by solving the eigenvalue problem). Then the Mohrs circle can be constructed as shown below:

The stress state for any rotation will be represented by a point either on one of the 3 circles, or in the shaded green area between the inner and outer circles.


MIT OpenCourseWare
http://ocw.mit.edu

### 2.080J / 1.573J Structural Mechanics

Fall 2013

For information about citing these materials or our Terms of Use, visit: http://ocw.mit.edu/terms.


[^0]:    ${ }^{1}$ Example 1.3 from Ugural and Fenster, Advanced Strength and Applied Elasticity, 2003

