Basic Physics of Underwater Acoustics

Reference used in this lecture: Lurton, X. 2002. An introduction to underwater acoustics. New York: Springer.

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Definitions

p: pressure, measured relative to hydrostatic, Pa

- ρ: density, measured *relative to hydrostatic*, kg/m³
- E: bulk modulus of the fluid, Pa, $\delta p = E [\delta \rho / \rho]$
- [u,v,w]: deflections in [x,y,z]-directions, relative to the hydrostatic condition, m

Then in one dimension *(pipe)* $p = E [-\delta u / \delta x]$



One-dimensional Case cont.



Let
$$p(x,t) = P_o \sin(\omega t - kx)$$



Insert this in the wave equation:

In water:

 $\rho \sim 1000 \text{ kg/m}^3$, E ~ 2.3e9 N/m² \rightarrow c ~ 1500 m/s

Wavelength $\lambda = 2\pi/k = 2\pi c/\omega = c/f$; **1kHz : 1.5m** *in water*

In Three Dimensions: A CUBE

Newton's Law:

 $p_{x} = -\rho \ u_{tt} \rightarrow p_{xx} = -\rho \ u_{ttx}$ $p_{y} = -\rho \ v_{tt} \rightarrow p_{yy} = -\rho \ v_{tty}$ $p_{z} = -\rho \ w_{tt} \rightarrow p_{zz} = -\rho \ w_{ttz}$

Constitutive Law:

- E
$$u_x = p / 3 \rightarrow - E u_{ttx} = p_{tt} / 3$$

- E
$$v_y = p / 3 \rightarrow - E v_{tty} = p_{tt} / 3$$

-
$$E w_z = p / 3 \rightarrow - E w_{ttz} = p_{tt} / 3$$

All directions deform uniformly



Lead to Helmholtz Equation:

$$p_{xx}+p_{yy}+p_{zz} = p_{tt} / c^2$$

or $\nabla^2 p = p_{tt} / c^2$

where ∇^2 is the LaPlacian operator

Particle Velocity

Consider one dimension again (Newton's Law): $p_x = -\rho \ u_{tt} \rightarrow p_x = -\rho \ (u_t)_t$

If $p(x,t) = P_o \sin(\omega t - kx)$ and $u_t(x,t) = U_{to} \sin(\omega t - kx) \rightarrow$

$$-kP_{o}\cos() = -\rho \omega U_{to}\cos() \rightarrow U_{to} = P_{o}/\rho c$$

Note velocity is in phase with pressure! [ρ c]: <u>characteristic impedance;</u> water: ρc ~ 1.5e6 Rayleighs "hard"

air:
$$\rho c \sim 500$$
 Rayleighs "soft"

In three dimensions:

 $\nabla p = -\rho V_t$ where

$$\nabla p = p_x i + p_y j + p_z k$$
 and
 $\underline{V} = u_t i + v_t j + w_t k$

Note equivalence of the following:

 $\lambda = c / f$ and $\omega / k = c$ There is <u>no</u> dispersion relation here; this is the only relationship between ω and k!

Consider Average Power through a 1D surface: $P(x) = [1/T] \int^{T} p(\tau,x) u_{t}(\tau,x) d\tau$ $= [1/T] \int^{T} P_{o} U_{to} \sin^{2}(\omega \tau - kx) d\tau$ $= P_{o} U_{to} / 2$ $= P_{o}^{2} / 2 \rho c = U_{to}^{2} \rho c / 2$ <u>Acoustic Intensity</u> in W/m²



Power per unit area is pressure times velocity

If impedance ρc is high, then it takes little power to create a given pressure level; but it takes a lot of power to create a given velocity level

Spreading in Three-Space

At time t_1 , perturbation is at radius r_1 ; at time t_2 , radius $r_2 \rightarrow \mathbf{P}(r_1) = \mathbf{P}_0^2(r_1) / 2 \rho c$ $\mathbf{P}(r_2) = \mathbf{P}_0^2(r_2) / 2 \rho c$

Assuming no losses in water; then $P(r_{2}) = P(r_{1}) r_{1}^{2} / r_{2}^{2} = P_{o}^{2} (r_{1}) r_{1}^{2} / 2 \rho c r_{2}^{2}$ and $P_{o}(r_{2}) = P_{o}(r_{1}) r_{1} / r_{2}$



Let $r_1 = 1$ meter (standard!) \rightarrow $P(r) = P_o^2(1m) / 2 \rho c r^2$ $P_o(r) = P_o(1m) / r$ $U_{to}(r) = P_o(1m) / \rho c r$

Pressure level and particle velocity decrease linearly with range

Decibels (dB)

10 * \log_{10} (ratio of two positive scalars):

Example: $x_1 = 31.6$; $x_2 = 1 \rightarrow 1.5$ orders of magnitude difference $10^* \log_{10}(x_1/x_2) = 15dB$ $10^* \log_{10}(x_2/x_1) = -15dB$

RECALL $\log(x_1^2/x_2^2) = \log(x_1/x_2) + \log(x_1/x_2) = 2 \log(x_1/x_2)$

In acoustics, acoustic intensity (power) is referenced to **1 W/m²** ; pressure is referenced to **1 μPa**

 $10^{*}\log_{10}[\mathbf{P}(\mathbf{r}) / 1 \text{ W/m}^{2}] = 10^{*}\log_{10}[\mathbf{P}_{o}^{2}(\mathbf{r}) / 2 \rho c] / 1 \text{ W/m}^{2}]$ $= 20^{*}\log_{10}[\mathbf{P}_{o}(\mathbf{r})] - 10^{*}\log_{10}(2\rho c)$ $= 20^{*}\log_{10}[\mathbf{P}_{o}(\mathbf{r}) / 1\mu\text{Pa}] - 120 - 65$

Spreading Losses with Range

Pressure level in dB at range r is $20 \log_{10} [P_o(r) / 1\mu Pa] - 185 =$ $20 \log_{10} [P_o(1m) / r / 1\mu Pa] - 185 =$ $20 \log_{10} [P_o(1m) / 1\mu Pa] -$ **20 \log_{10} [r]**-185

Example: At 100m range, we have lost 40dB or *four orders of magnitude* in sound intensity 40dB or *two orders of magnitude* in pressure (and particle velocity)

Attenuation Losses with Range

Acoustic power <u>does</u> have losses with transmission distance – primarily related to relaxation of boric acid and magnesium sulfate molecules in seawater. Also bubbles, etc.

At 100 Hz, ~1dB/1000km: OK for thousands of km, ocean-scale seismics and communications

At 10kHz, ~1dB/km: OK for ~1-10km, long-baseline acoustics

At 1MHz, 3dB/10m: OK for ~10-100m, imaging sonars, Doppler velocity loggers



The Piezo-Electric Actuator

strain = constant X electric field $\varepsilon = \mathbf{d} \times \mathbf{E}$ or $\Delta t/t = dX(V/t)$ where d = 40-750 x 10^{-12} m / V \rightarrow Drive at 100V, we get only 4-75 nm thickness change!





The Piezo-Electric Sensor

electric field = constant X stress $E = g X \sigma$ or $V = t g \sigma$ where g = 15-30 x 10⁻³ Vm/N

<u>Ideal Actuator</u>: Assume the water does not impede the driven motion of the material <u>Ideal Sensor</u>: Assume the sensor does not deform in response to the water pressure waves Typical Transducer:
 $120 to 150 dB re 1\mu Pa, 1m, 1V$ means
or
or
1.30 Pa at 1m for each Volt appliedTypical Hydrophone:
 $-220 to -190 dB re 1\mu Pa, 1V$ means
or 10^{-11} to $10^{-9.5}$ V for each μ Pa incidentor

10⁻⁵ to 10^{-3.5} V for each Pa incident

So considering a transducer with 16Pa at 1m per Volt, and a hydrophone with 10⁻⁴ V per Pa: If V = 200V, we generate 3200Pa at 1m, or 3.2Pa at 1km, assuming spreading losses only; The hydrophone signal at this pressure level will be 0.00032V or 320μV ! 2.017J Design of Electromechanical Robotic Systems Fall 2009

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