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### 3.23 Electrical, Optical, and Magnetic Properties of Materials

Fall 2007

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### 3.23 Fall 2007 - Lecture 2

## More practical info

- Problem sets - out on Wed (and posted on Stellar), due by 5 pm of the following weekerfd (after that 75\%, after Thu 5pm 50\%, after Fri 5pm 25\% )
- ~11 in total, $30 \%$ of the grade
- Sometimes I mention homework - it's not the "Problem Set" @ Poilvert, Bonnet


## Homework

- Take notes
- Revise posted lecture
- Study posted or assigned material (TEXTBOOKS - do you have them ?)
- Meet with TAs or Instructor:

Marzari Office Hours - Monday 4-5 pm
Poilvert Office Hours - Tuesday 4-5pm

## Last time: Wave mechanics

1. Particles, fields, and forces
2. Dynamics - from Newton to Schroedinger
3. De Broglie relation $\lambda \bullet p=h$
4. Waves and plane waves
5. Harmonic oscillator

## Time-dependent Schrödinger's equation <br> (Newton's $2^{\text {nd }}$ law for quantum objects)



1925-onwards: E. Schrödinger (wave equation), W. Heisenberg (matrix formulation), P.A.M. Dirac (relativistic)

## Plane waves as free particles

Our free particle $\Psi(\vec{r}, t)=A \exp [i(\vec{k} \cdot \vec{r}-\omega t)]$ satisfies the wave equation:

$$
-\frac{\hbar^{2}}{2 m} \nabla^{2} \Psi(\vec{r}, t)=i \hbar \frac{\partial \Psi(\vec{r}, t)}{\partial t} \quad \text { (provided } \quad E=\hbar \omega=\frac{p^{2}}{2 m}=\frac{\hbar^{2} k^{2}}{2 m} \text { ) }
$$

Stationary Schrödinger's Equation (I)

$$
\begin{gathered}
-\frac{\hbar^{2}}{2 m} \nabla^{2} \Psi(\vec{r}, t)+V(\vec{r}, *) \Psi(\vec{r}, t)=i \hbar \frac{\partial \Psi(\vec{r}, t)}{\partial t} \\
\Psi(\vec{r}, t)=\text { ANSATE } \varphi(\vec{r}) f(t) \\
-\frac{\hbar^{2}}{2 m} \nabla^{2}(\varphi f)+V(\vec{r}) \varphi f=i \hbar \frac{\partial(\varphi t)}{\partial t} \\
-\frac{\hbar^{2}}{2 m} f \nabla^{2} \varphi+V \varphi f=i \hbar \varphi \frac{\partial f}{\partial t} / \varphi f
\end{gathered}
$$

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$$
\begin{aligned}
& \underbrace{-\frac{\hbar^{2} \sigma^{2} \varphi}{2 m}+V}_{\vec{r}}=\underbrace{f \frac{1}{\varphi} \frac{\partial f}{\partial t}}_{f}=\cdots \\
& =\operatorname{ConsiAAT}=E \\
& -\frac{\hbar^{2}}{2 m} \frac{\nabla^{2} \varphi}{\varphi}+V=E \quad i \hbar f^{\prime} \frac{\partial f}{f}=E
\end{aligned}
$$

## Stationary Schrödinger's Equation (II)


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## Stationary Schrödinger's Equation (III)

$$
\left[-\frac{\hbar^{2}}{2 m} \nabla^{2}+V(\vec{r})\right] \varphi(\vec{r})=E \varphi(\vec{r})
$$

1. It's not proven - it's postulated, and it is confirmed experimentally
2. It's an "eigenvalue" equation: it has a solution only for certain values (discrete, or continuum intervals) of $E$
3. For those eigenvalues, the solution ("eigenstate", or "eigenfunction") is the complete descriptor of the electron in its equilibrium ground state, in a potenitial $\mathrm{V}(\mathrm{r})$.
4. As with all differential equations, boundary conditions must be specified
5. Square modulus of the wavefunction = probability of finding an electron

Free particle: $\Psi(x, t)=\phi(x) f(t)$

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Infinite Square Well (II)

$$
\begin{aligned}
& \left(\left.(\sin k x)\right|_{x=a}=0\right. \\
& C \sin k a=0 \\
& k a=n \pi \quad n=0,+1,-2, \cdots
\end{aligned}
$$

Infinite Square Well (III)


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## The power of carrots

- $\beta$-carotene



## Physical Observables from Wavefunctions

- Eigenvalue equation:

$$
\left[-\frac{\hbar^{2}}{2 m} \frac{d^{2}}{d x^{2}}+V(x)\right] \varphi(x)=E \varphi(x)
$$

- Expectation values for the operator (energy)
$\lambda \mid W$

$$
E=\int \varphi^{*}(x)\left[-\frac{\hbar^{2}}{2 m} \frac{\partial^{2}}{\partial x^{2}}+V(x)\right] \varphi(x) d x \quad E=\frac{h^{2}}{\delta m}\left(\frac{n^{2}}{a^{2}}\right)
$$

## Particle in a 2-dim box

$$
\begin{aligned}
& -\frac{\hbar^{2}}{2 m}\left(\frac{\partial^{2}}{\partial x^{2}}+\frac{\partial^{2}}{\partial y^{2}}\right) \varphi(x, y)=E \varphi(x, y) \\
& \varphi(x, y)=X(x) Y(y) \\
& -\frac{\hbar^{2}}{2 m} y \frac{\partial^{2} x}{\partial x}-\frac{\hbar^{2}}{2 m} \times \frac{\partial^{2} y}{\partial y^{2}}=E X y \\
& - \\
& \frac{\hbar^{2}}{2 m} \frac{1}{2 m} \frac{\partial^{2} x}{\partial n_{323}^{2}}=E=-\frac{\hbar^{2}}{2 m} \frac{1}{\lambda^{2}} \frac{\partial^{2} y}{\partial y^{2}}
\end{aligned}
$$

## Particle in a 2-dim box

$$
\varphi(x, y)=C \sin \left(\frac{l \pi x}{a}\right) \sin \left(\frac{m \pi y}{b}\right)
$$



$$
E=\frac{h^{2}}{8 m}\left(\frac{l^{2}}{a^{2}}+\frac{m^{2}}{b^{2}}\right)
$$


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$$
\begin{gathered}
\text { Particle in a 3-dim box } \\
-\frac{\hbar^{2}}{2 m}\left(\frac{\partial^{2}}{\partial x^{2}}+\frac{\partial^{2}}{\partial y^{2}}+\frac{\partial^{2}}{\partial z^{2}}\right) \varphi(x, y, z)=E \varphi(x, y, z)
\end{gathered}
$$

## Particle in a 3-dim box: Farbe defect in halides ( $e^{-}$bound to a negative ion vacancy)



# From Carl Zeiss to MIT... 

Text removed due to copyright restrictions. Please see
Avakian, P., and Smakula, A. "Color Centers in Cesium Halide Single Crystals."
Physical Review 120 (December 1960): 2007.

## Light absorption/emission



Courtesy M. Bawendi and Felice Frankel.
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## Metal Surfaces (I)

$\left[-\frac{\hbar^{2}}{2 m} \frac{d^{2}}{d x^{2}}+V(x)\right] \varphi(x)=E \varphi(x)$


## Metal Surfaces (II)



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## Scanning Tunnelling Microscopy



Figure by MIT OpenCourseWare.

## Scanning Tunnelling Microscopy, cont.



$$
\begin{aligned}
& \mathrm{I} / \mathrm{V} \propto \rho \mathrm{e}^{-2 \kappa s} \\
& \kappa=\left(\frac{2 \mathrm{~m} \phi}{\hbar^{2}}\right)^{1 / 2}=1.1 \AA^{-1} \\
& \rho=\text { density of states }
\end{aligned}
$$

## Wavepacket tunnelling through a nanotube


http://newton.phy.bme.hu/education/schrd/index.html
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http://www.quantum-physics.polytechnique.fr

