# 12.815 Atmospheric Radiation Fall 2008

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### Energy Levels in Molecules, contd.

e.g. energy level diagram for O2



 $\label{eq:conversions: 1ev = 8067 cm^{-1} = 1.24 \ \mu\text{m} \ (\text{near infrared}) \\ 2ev = 16134 \ \text{cm}^{-1} = 0.62 \ \mu\text{m} \\ 3ev = 24201 \ \text{cm}^{-1} = 0.41 \ \mu\text{m} \\ 5ev = 40335 \ \text{cm}^{-1} = 0.248 \ \mu\text{m} \ (\text{ultraviolet}) \\ \end{array}$ 

Absorption and emission by gases

For transitions between two states, initial i and final f, the Einstein "b" coefficient is given by

$$b_{fi} = \frac{1}{4\pi} \times \frac{8\pi^{3} v_{fi}}{3 h c} \left| \int \Phi_{f}^{*} \left( \vec{\mu} \ \Phi_{i} \right) dV \right|^{2} \left| \int \sum_{f}^{*} \sum_{i} d\sigma \right|^{2}$$

and we need this to be  $\gg$  0 for strong ("allowed") absorption or emission. This defines "selection rules".

 $\Phi(\mathbf{r}, \theta, \phi)$  = space-dependent part of the wavefunction describing state of molecule.

 $\sum(\sigma) = \text{spin-dependent part of wavefunction} \\ (\text{total wavefunction} = \Phi . \Sigma )$ 

 $\vec{\mu}$  = dipole moment operator =  $\vec{er}$ 

(a) Selection rules from  $\mu$  (dipole)

For strong absorption/emission need  $\mu \neq 0$ 

- (i) <u>pure rotational transitions</u> permanent dipole <u>required</u>
- (ii) <u>vibrational transitions</u> where vibration does <u>not</u> produce a dipole moment permanent dipole <u>required</u>
- (iii) <u>vibrational transitions</u> where vibration produces a (transient) dipole moment permanent dipole <u>not required</u>.
- (iv) when vibrational transitions allowed then simultaneous vibrational + rotational transitions allowed
- (v) <u>electronic transitions</u> permanent dipole <u>not required</u> (always a dipole between nuclei and electrons)

(b) Selection rules from  $\Phi_i$  and  $\Phi_f$  (symmetry)

Since  $\bar{\mu}$  (or  $\bar{r}$ ) is an <u>odd</u> function in space, and we require total integrated  $\Phi_{f}^{*} \bar{\mu} \Phi_{i}$  to be an <u>even</u> function to prevent significant cancellation between volume elements, then  $\Phi_{f}^{*}$  and  $\Phi_{i}$  <u>must have opposite parity</u>.

(La Porte's Rule). Formally this leads to:

i) for rotational energy transitions in a linear molecule

 $\Delta J = \pm 1$ 

(<u>special case</u> is  $\Delta J = 0, \pm 1$  for a non-diatomic simultaneous rotational-vibrational transition where vibration is  $\perp$  to axis)

ii) for vibrational energy transitions in all molecules

$$\Delta v = \pm 1 \qquad (i.e. \text{ generally } v=0 \rightarrow v=1)$$
(i) (f)

 $\frac{(i.e. \text{ overtones and } \underline{\text{combinations}} \text{ are } \underline{\text{not}} \text{ allowed})}{(n_{v})} \frac{(v_{1} + v_{2})}{(v_{1} + v_{2})}$ 

iii) for electronic transitions in linear molecules The symbols provide the needed information on parity of wave function  $\begin{pmatrix} 2S+1 \\ \downarrow \\ u, g \end{pmatrix}$ i.e.  $g \leftrightarrow u$  $+ \Leftrightarrow \begin{pmatrix} allowed transitions \end{pmatrix}$ 

(c) selection rules from 
$$\sum$$
 (spin)

For  $b_{ij} \neq 0$  need  $\int \sum_{i=1}^{s} \sum_{j=1}^{s} d\sigma \neq 0$  but since  $\sum_{i=1}^{s} \sum_{j=1}^{s} \sigma_{ij} = \sum_{j=1}^{s} \sum_{j=1}^{s} \sigma_{ij}$ 

Hence, transitions cannot involve a change of spin

This is important for electronic transitions. Note that spin is also given in the state symbol for electronic states.  $\begin{pmatrix} 2S+1 \\ u, g \end{pmatrix}$ 

e.g. O2 (see Energy Level diagram)

$${}^{3}\sum_{g}^{-} \xrightarrow{\text{allowed}} {}^{3}\sum_{u}^{+} \text{ (Herzberg bands)}$$

$$\xrightarrow{\text{allowed}} {}^{3}\sum_{u}^{-} \text{ (Schumann-Runge bands)}$$

$$\xrightarrow{\text{forbidden*}} {}^{1}\sum_{g}^{+} \text{ (red "atmospheric bands)}$$

$$\xrightarrow{\text{forbidden*}} {}^{1}\Delta_{g} \text{ (infrared "atmospheric" bands)}$$

(\* observed as very weak "magnetic dipole" transitions)

Shape of absorption and emission lines

(a) "<u>natural</u>" broadening due to (Heisenberg\*) uncertainty in energy of upper and lower levels. Integrated line strength =  $S_{i0} = \int_{0}^{\infty} k_v \, dv$ 



(b) "<u>doppler</u>" broadening due to motion toward and away from observer (use Maxwellian distribution of velocities)

$$S_{i0} = \int_{0}^{\infty} k_{v} dv = \int_{0}^{\infty} \frac{S_{i0}}{\sqrt{\sqrt{\pi}}} \exp\left[-\left(\frac{v - v_{0}}{\gamma}\right)^{2}\right] dv$$



(c) "<u>pressure</u>" broadening due to collisional and long distance interactions perturbing the energy levels

$$S_{i0} = \int k_{v} dv = \int_{0}^{\infty} \left[ \frac{k_{0}}{1 + \left( \frac{v - v_{0}}{\Gamma/4\pi} \right)^{2}} \right] dv$$

$$\Gamma \propto \text{ collisional frequency} \checkmark$$

$$\sum_{i=1}^{\infty} \left[ \frac{k_{0}}{1 + \left( \frac{v - v_{0}}{\Gamma/4\pi} \right)^{2}} \right] dv$$

$$\sum_{i=1}^{\infty} \left[ \frac{NOTE}{Collision freq.(sec^{-1}) \text{ in a } 2-component gas (m_{1}, r_{1}, n_{1} \text{ and } m_{2}, r_{2}, n_{2}) \text{ is:} \right]$$

$$\sum_{i=1}^{\infty} \left[ \frac{k_{0}}{1 + \left( \frac{v - v_{0}}{\Gamma/4\pi} \right)^{2}} \right] dv$$

i.e. 
$$\Gamma = \Gamma_0$$
 (1 atm, 273K)  $P^{\overline{\alpha}} T^{-\overline{\beta}} = \left(\overline{\alpha} \approx 1, \ \overline{\beta} \approx \frac{1}{2}\right) (P = nKT)$ 

 $(\Gamma_0, \overline{\alpha}, \overline{\beta} \text{ from laboratory measurements})$ 

$$(k_0 = 4S_{i0}/\Gamma)$$



(d) Summary





#### Absorption bands

Can get rotational, vibrational, and electronic transitions together:

(a) rotational bands (e.g.  $H_2O$ ) – due to  $\Delta J = \pm 1$ 



#### (b) vibrational-rotation bands – due to $\Delta J = 0, \pm 1$ and $\Delta v = \pm 1$





## (c) electronic – vibration – rotation bands



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