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12.510 Introduction to Seismology Spring 2008

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 $\diamond$  Go back to the observations again and look at the deviation from the model.

**Objective**: we need to find a model that minimize  $\delta t$ 

 $T_{obs}$ = 3D structure + error in source location + noise (instrument, measurement)

Hypocenter: to,x,y,z (or 
$$\theta,\phi,r$$
)

 $\Rightarrow \delta t = \delta t_{3D} + \delta t_{mislocation} + \delta t_{noise}$ 

Errors caused by the model and the source location are usually combined together, and the noise is assumed to be white, with a Gaussian distribution.

If we know the 1D model, we can apply **Snell's law** to estimate the geometry:

$$T_{obs} = \int_{\substack{3D\\ Ray\\ Path}} \frac{1}{c(\underline{x})} dl$$

Note that

$$\nabla T(x) = \frac{1}{c(x)} \underline{k}$$

If c change, the ray path changes. We end up with a nonlinear problem. Thus, we should try to linearize the inversion, using **the Fermat's Principle**.





If we change a little bit the ray path around the optimum, we'll end up with a small change in travel time. We have two kinds of "deviations" from the reference ray:

1. First contribution: effect of changes in velocity  $\Delta c$ .

2. Second contribution: effect of change in the ray path. Fermat's principle says that we can ignore it.

## Linearization of the travel time

## Travel time residual:

$$Observation = \delta t_{3D} = T_{obs} - T_{ref} = \int_{\substack{true \\ 3D \\ structure}} \frac{1}{c(\underline{x})} dl - \int_{\substack{reference \\ 3D \\ path}} \frac{1}{c_0(\underline{x})} dl_0$$

$$Fermat's \text{ Princeple} \approx \int_{\substack{reference \\ 3D \\ path}} \frac{1}{c(\underline{x})} dl_0 - \int_{\substack{reference \\ 3D \\ path}} \frac{1}{c_0(\underline{x})} dl_0$$

$$= \int_{\substack{reference \\ 3D \\ path}} \frac{\Delta c}{c_0^2} dl_0 = \int_{\substack{reference \\ 3D \\ path}} (s(\underline{x}) - s_0(\underline{x})) dl_0$$

$$= \int_{\substack{reference \\ 3D \\ path}} (\Delta s(\underline{x})) dl_0$$

In linearizing the problem, we get rid of the unknown ray. We can do our calculation in a reference earth model.

- $\diamond$  The travel time tomography is an iterative process:
- ♦ Create 1D model
- $\diamond$  Ray tracing and get new rays in the model
- $\diamond$  Update ray geometry
- $\diamond$  Get the reference ray related to the 3D

$$\delta t = T_{obs} - T_{ref} (3D)$$

(The reference model does not have to be a 1D model.)

## Linearization of the hypocenter mislocation

$$\begin{split} \delta t_{misloc} &= \frac{\partial T}{\partial t_o} \delta t_o + \frac{\partial T}{\partial x} \delta x + \frac{\partial T}{\partial y} \delta y + \frac{\partial T}{\partial z} \delta z \\ \delta &\otimes = \frac{\partial \otimes}{\partial \oplus} \delta \oplus \end{split}$$

with  $t_0$  the origin time, (x, y, z) the location of the earthquake.

Then we try to solve for  $\Delta s$ ,  $\delta t$ ,  $\delta x$ ,  $\delta y$ ,  $\delta z$ .

## Inverse Problem



First, we need to discretize the problem, i.e. **parameterization**:

$$\mu(\underline{r}) = \sum_{k=1}^{M} \gamma_k h_k(\underline{r})$$

 $\gamma_k$ : weight  $h_k$ : basis function

For example, plane wave summation:

$$\phi = \iiint \Phi(x, y, z, \omega) e^{i(\underline{k} \cdot \underline{x} - \omega t)} d\underline{k} d\omega$$

Then, we inject  $\mu$  into the equation:

$$d_i = \sum_{k=1}^{M} \left\{ \int G_i(\underline{r}) h_k(\underline{r}) d\underline{r} \right\} \gamma_k$$

kernel projected on the basis function h

$$d_i = \sum_{k=1}^{M} A_{ik} \gamma_k$$
 for i=1,2,..., N data  
sensitivity matrix kernel projected on the basis functions

projector  

$$\underline{d} = \underline{\underline{Am}}$$
  $m = (\gamma_1, \gamma_2, ..., \gamma_M)$   
model (includes the values of the weight)  
data

where

 $G_i$ : Green functions, solution of a point source. We need to do a convolution with a point perturbation in order to get the observations.

M: Number of model parameters.

N: Number of observations.

In general,  $M \neq N$ . As a consequence, the matrix <u>A</u> is not square.



M columns

Multiplying the equation by  $A^{T}$ , we can get the solution:

$$\rightarrow \underline{\hat{m}} = (A^T A)^{-1} A^T \underline{d}$$
 Generalized Least Squares inversion  
$$\|\hat{m} - (A^T A)^{-1} A^T \underline{d}\| = \varepsilon \rightarrow \text{Goal is to minimize } \varepsilon.$$

Back to our specific inverse problem:

$$\delta t = \int \Delta s.dl \rightarrow \text{parametrize } \Delta s$$
$$\Delta s = \sum_{k=1}^{M} \gamma_k h_k$$

One way is to take  $h_k$  as a series of cells/blocks, with a value for  $\underline{x}$  inside the cell k and zero otherwise. We have

$$\delta t_{i} = \int \Delta s dl = \sum_{k=1}^{M} \Delta s_{k} (dl)_{ik}$$

where

*i* : event-station pair.

 $(dl)_{ik}$ : path length.

Rewrite the equation in the matrix form:

$$A_{ik} \cdot m = \begin{bmatrix} \Delta l_{11} & \dots & \Delta l_{1M} \\ \vdots & \vdots & \vdots \\ \Delta l_{N1} & \dots & \Delta l_{NM} \end{bmatrix} \begin{bmatrix} \Delta s_1 \\ \vdots \\ \Delta s_M \end{bmatrix} = \begin{bmatrix} \delta t_1 \\ \vdots \\ \delta t_N \end{bmatrix}$$

where each ray gives a row in the matrix.

We have an average wavespeed along the ray. In order to construct a model vector, we need to get data from different rays crossing each other.

A is a sparse matrix. If we look at one ray:



A will have only ~100 elements non-zero. The good thing about sparse matrix is that  $A^T A$  is approximately diagonal. The problem is that there are many singularities, which make the inversion unstable (in that case, we need to add a damping factor or regularize the problem). One possibility is to not use cells of the same size. Consequently, it reduces the number of cells; the inverse matrix is less singular. Nevertheless, the computation time increases.

Another way is take  $h_k$  as spherical harmonics (in global seismology).