Supplementary Notes for Chapter 5 The Calculus of Thermodynamics

#### **Objectives of Chapter 5**

- 1. to understand the framework of the Fundamental Equation including the geometric and mathematical relationships among derived properties (*U*, *S*, *H*, *A*, and *G*)
- 2. to describe methods of derivative manipulation that are useful for computing changes in derived property values using measurable, experimentally accessible properties like *T*, *P*, *V*,  $N_i$ ,  $x_i$ , and  $\rho$ .
- 3. to introduce the use of Legendre Transformations as a way of alternating the Fundamental Equation without losing information content

# Starting with the combined 1<sup>st</sup> and 2<sup>nd</sup> Laws and Euler's theorem we can generate the Fundamental Equation:

## Recall for the combined 1<sup>st</sup> and 2<sup>nd</sup> Laws:

- Reversible, quasi-static
- Only *PdV* work
- Simple, open system (no KE, PE effects)
- For an *n* component system

$$d\underline{U} = Td\underline{S} - Pd\underline{V} + \sum_{i=1}^{n} (H - TS)_i \ dN_i$$

$$d\underline{U} = Td\underline{S} - Pd\underline{V} + \sum_{i=1}^{n} \mu_i dN_i$$

#### and Euler's Theorem:

• Applies to all smoothly-varying homogeneous functions *f*,

f(a,b,...,x,y,...)

where  $a, b, \ldots$  intensive variables are homogenous to zero order in mass and x, y, extensive variables are homogeneous to the 1<sup>st</sup> degree in mass or moles (N).

• *df* is an exact differential (not path dependent) and can be integrated directly

if 
$$Y = ky$$
 and  $X = kx$  then

$$f(a,b, ..., X, Y, ...) = k f(a,b, ..., x, y, ...)$$

and

$$x\left(\frac{\partial f}{\partial x}\right)_{a,b,\dots,y,\dots} + y\left(\frac{\partial f}{\partial y}\right)_{a,b,\dots,x,\dots} + \dots = (1)f(a,b,\dots,x,y,\dots)$$

#### **Fundamental Equation:**

- Can be obtained via Euler integration of combined 1<sup>st</sup> and 2<sup>nd</sup> Laws
- Expressed in Energy ( $\underline{U}$ ) or Entropy ( $\underline{S}$ ) representation

$$\underline{U} = f_u [\underline{S}, \underline{V}, N_1, N_2, ..., N_n] = T \underline{S} - P \underline{V} + \sum_{i=1}^n \mu_i N_i$$

or

$$\underline{S} = f_s [\underline{U}, \underline{V}, N_1, N_2, ..., N_n] = \frac{\underline{U}}{T} + \frac{P}{T} \underline{V} - \sum_{i=1}^n \frac{\mu_i}{T} N_i$$

# The following section summarizes a number of useful techniques for manipulating thermodynamic derivative relationships

Consider a general function of n + 2 variables

$$F(x, y, z_3, ..., z_{n+2})$$

where  $x \equiv z_1$ ,  $y \equiv z_2$ . Then expanding via the rules of multivariable calculus:

$$dF = \sum_{i=1}^{n+2} \left( \frac{\partial F}{\partial z_i} \right) dz_i$$

Now consider a process occurring at constant F with  $z_3, ..., z_{n+2}$  all held constant. Then

$$dF = 0 = \left(\frac{\partial F}{\partial x}\right)_{y, z_3, \dots} dx + \left(\frac{\partial F}{\partial y}\right)_{x, z_3, \dots} dy$$

Rearranging, we get:

#### **Triple product** *"x-y-z-(-1)* **rule" for** *F*(*x*,*y*)**:**

$$\left(\partial F / \partial x\right)_{y} \left(\partial x / \partial y\right)_{F} \left(\partial y / \partial F\right)_{x} = -1$$

example:

$$(\partial H / \partial T)_P (\partial T / \partial P)_H (\partial P / \partial H)_T = -1$$

### Add another variable to F(x,y):

$$\left(\partial F / \partial y\right)_{x} = \frac{\left(\partial F / \partial \phi\right)_{x}}{\left(\partial y / \partial \phi\right)_{x}}$$

example: F(x, y) = S(P, H) and  $\phi = T$  then

$$\left(\frac{\partial S}{\partial H}\right)_{P} = \frac{\left(\frac{\partial S}{\partial T}\right)_{P}}{\left(\frac{\partial H}{\partial T}\right)_{P}} = \frac{C_{p}/T}{C_{p}} = 1/T$$

#### **Derivative inversion for** F(x,y)**:**

$$(\partial F / \partial y)_x = 1/(\partial y / \partial F)_x$$

example:

$$(\partial T / \partial S)_P = 1/(\partial S / \partial T)_P = T / C_p$$

### Maxwell's reciprocity theorem:

Applies to all homogeneous functions, e.g. F(x,y, ..)

$$\left\lfloor \frac{\partial (\partial F / \partial x)_{y,\dots}}{dy} \right\rfloor_{x,\dots} = \left\lfloor \frac{\partial (\partial F / \partial y)_{x,\dots}}{\partial x} \right\rfloor_{y,\dots} \text{ or } F_{xy} = F_{yx}$$

example:

$$d\underline{U} = Td\underline{S} - Pd\underline{V} + \sum_{i=1}^{n} \mu_i dN_i$$
$$(\partial T / \partial \underline{V})_{\underline{S},N} = \underline{U}_{\underline{S}\underline{V}} = \underline{U}_{SV} = -(\partial P / \partial \underline{S})_{\underline{V},N} = \underline{U}_{\underline{V}\underline{S}} = \underline{U}_{VS}$$

Legendre Transforms:

$$\begin{array}{c} (x_i, \xi_i) \\ (\underline{S}, T) \\ (\underline{V}, -P) \\ (N_i, \mu_i) \\ (\underline{x}_i, F_i) \\ (\underline{a}, \sigma) \end{array} \right\}$$
 Conjugate coordinates (extensive, intensive)

## General relationship

$$y^{(0)} = f[x_1, ..., x_m]$$
 (basis function)

$$y^{(k)} = y^{(0)} - \sum_{i=1}^{k} \xi_i x_i \quad (k^{th} \ transform)$$

 $\underline{U} = f[\underline{S}, \underline{V}, N_1, \dots N_n]$ 

 $y^{(1)} = \underline{A} = \underline{U} - T\underline{S}$ 

or by changing variable order to  $U = f(\underline{V}, \underline{S}, N_1, ..., N_n),$ 

$$y^{(1)} = \underline{H} = \underline{U} + P\underline{V}$$

$$dy^{(k)} = -\sum_{i=1}^{k} x_i d\xi_i + \sum_{i=k+1}^{m} \xi_i dx_i$$

$$dy^{(1)} \equiv d\underline{A} = -\underline{S}dT - Pd\underline{V} + \sum_{i=1}^{n} \mu_i dN_i$$
  
or

$$dy^{(1)} \equiv d\underline{H} = Td\underline{S} + \underline{V}dP + \sum_{i=1}^{n} \mu_i dN_i$$

$$y^{(m)} = y^{(0)} - \sum_{i=1}^{m} \xi_i x_i = 0$$

$$y^{(n+2)} = 0$$
 (total transform with  $m = n+2$ )

$$dy^{(m)} = -\sum_{i=1}^{m} x_i d\xi_i = 0$$

$$dy^{(n+2)} = -\underline{S}dT + \underline{V}dP - \sum_{i=1}^{n} N_i d\mu_i = 0$$
(Gibbs-Duhem Equation)

# **Relationships among Partial Derivatives of Legendre Transforms**

$$y_{ij}^{(k)} = y_{ji}^{(k)} = \frac{\partial^2 y^{(k)}}{\partial x_i \partial x_j} \qquad (Maxwell relation) \qquad \xi_i \equiv y_i^{(0)} = \left(\frac{\partial y^{(0)}}{\partial x_i}\right)_{x_j[i]}$$
$$y_{1i}^{(0)} = \frac{\partial^2 y^{(0)}}{\partial x_1 \partial x_i} \qquad y_{11}^{(0)} = \frac{\partial^2 y^{(0)}}{\partial x_1^2} \qquad y_i^{(1)} = \begin{cases} -x_i & i = 1 \\ \xi_i & i > 1 \end{cases}$$
$$[NB: \quad \xi_i = y_i^{(0)} \text{ as well for } i > 1]$$

#### **Reordering and Use of Tables 5.3-5.5**

Table 5.3 – 2<sup>nd</sup> & 3<sup>rd</sup> order derivatives of  $[y_{ij}^{(1)} \text{ and } y_{ijk}^{(1)}]$  in terms of  $y_{ii}^{(0)}$ , etc

- Table 5.4 Relations between  $2^{nd}$  order derivatives of  $j^{ih}$  Legendre transform  $y_{ik}^{(j)}$  and the basis function  $y_{ik}^{(0)}$
- Table 5.5 Relationships among 2<sup>nd</sup> order derivatives of  $j^{th}$  Legendre transform  $y_{ik}^{(j)}$  to (*j-q*) transform  $y_{ik}^{(j-q)}$

<sup>\10.40\</sup>Ch5-Calc. of Thermo. Suppl. Notes