

6.581/20.482J Problem Set #2

Due 5:00 p.m. Tuesday, March 14th, 2006

1. Modeling the dynamics of a trapped ion. Here, we will explore the implementation of several integration schemes in molecular dynamics. Figure 1 shows the system of interest—a single mobile ion trapped by a square of like-charged ions.

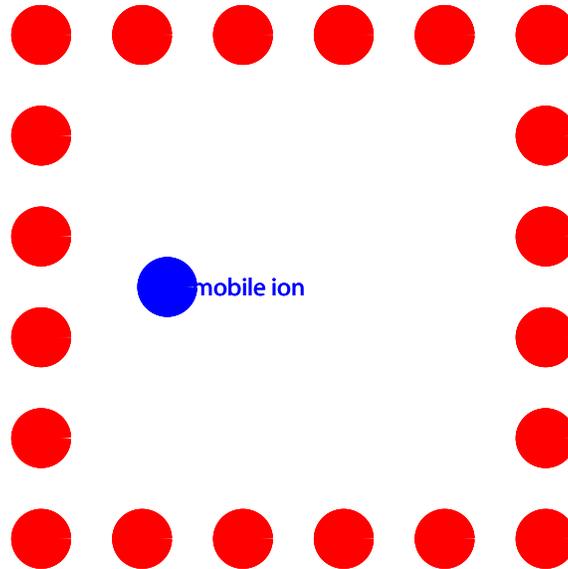


Figure 1: Physical system for molecular dynamics simulation.

The MATLAB script `mol_dyn.m` contains the description of the system, as well as the force and energy functions for both electrostatic interactions:

$$E_{ij}^{\text{elec}} = \frac{q_i q_j}{4\pi\epsilon_0 r_{ij}} \quad (1)$$

and van der Waals interactions:

$$E_{ij}^{\text{vdw}} = 4\sqrt{\epsilon_i \epsilon_j} \left[ \left( \frac{\sqrt{\sigma_i \sigma_j}}{r_{ij}} \right)^{12} - \left( \frac{\sqrt{\sigma_i \sigma_j}}{r_{ij}} \right)^6 \right] \quad (2)$$

The van der Waals parameters  $\epsilon$  and  $\sigma$  describe the energy of the most favorable interaction ( $\epsilon$ ) and the distance at which the interaction becomes unfavorable ( $\sigma$ ). The parameters for three different ions are displayed below.

Ion	m (g/mol)	q (e)	$\sigma$ (Å)	$\epsilon$ (kcal/mol)
F <sup>-</sup>	18.998	-1.0	2.73295	0.72000
Cl <sup>-</sup>	35.453	-1.0	4.41724	0.11779
Br <sup>-</sup>	79.904	-1.0	4.62376	0.09000

Please note that in molecular level simulations, distance is measured in Å ( $1 \times 10^{-10}$  m), charges in units of proton charge (e), masses in units of g/mol, and energies in kcal/mol; in these units,  $\epsilon_0$  is  $2.397 \times 10^{-4} \text{ e}^2 \cdot \text{Å} \cdot (\text{kcal/mol})^{-1}$ . Although this choice of units leads to a unit of time equal to  $4.888821 \times 10^{-14}$  s, a conversion to a ps time unit is generally implemented (as in the MATLAB script `mol_dyn.m`).

- (a) Implement the Forward Euler integration scheme, and propagate the dynamics of the ion with a 0.01 ps timestep.
  - (b) Describe the dynamic trajectory of the mobile ion.
  - (c) What happens to the energy of the system (kinetic, potential, and total) as a function of time?
  - (d) Repeat the simulation with a 0.001 ps time step, and repeat (b) and (c), and comment on the differences.
  - (e) Implement the Velocity Verlet integration scheme, and again propagate the dynamics of the ion with a 0.01 ps timestep. Repeat (b) and (c) for this trajectory. How could you test the validity of this trajectory? Verify your idea.
  - (f) Repeat (d) using a 0.1 ps time step, and comment on the differences.
  - (g) Repeat the simulation using both  $\text{Cl}^-$  and  $\text{Br}^-$  as the mobile ion (choose the integration method you feel is most appropriate). Compare the trajectories with the results from the  $\text{F}^-$  simulation. How does the trajectory of the ion change with ion type? How do the energetics of the ion change?
  - (h) Repeat the simulation with an initial position of  $-9.0 \text{ \AA}$ , again for each ion type. Discuss your observations of the trajectory and the energetics.
2. **Long-range electrostatic interactions.** In this problem, we will use eigendecomposition, via MATLAB's `eig` function, to illustrate that sometimes simple approximations can significantly reduce computational requirements (memory and time) while not sacrificing much accuracy. Recall that for a matrix  $A$ , an eigenvalue  $\lambda$  and a corresponding eigenvector  $x$  satisfy

$$Ax = \lambda x, \quad (3)$$

and that a symmetric matrix has a complete set of orthogonal eigenvectors.

Figure 2 shows two clusters of point charges. Each atom has a partial charge on it, which we model as a point charge at the atom center. Let the set of atomic charges in cluster 1 be denoted by the vector  $q$ ; the set in cluster 2, by  $s$ . The charges  $q$  produce a potential field in cluster 2; let the vector of potentials at the atom centers in cluster 2 be  $\phi$ . This vector of potentials produced by  $q$  can be found by the matrix multiplication

$$\phi = Pq \quad (4)$$

where  $P$  is the potential matrix; the  $i^{\text{th}}$  column of  $P$  is the set of potentials produced at the atom centers in cluster 2, due to a unit charge at point  $i$  in cluster 1. (Thus an arbitrary charge distribution  $q$  in cluster 1 produces a scaled sum of the responses due to each charge separately.) We have set up a symmetric physical system, so  $P$  will be symmetric. In general  $P$  is unsymmetric, so the singular value decomposition (SVD) is used for analysis rather than eigendecomposition. Trefethen and Bau's Numerical Linear Algebra contains a very good exposition of the SVD.



Figure 2: Electrostatic system: two well-separated, symmetric clusters of charges.

If the clusters are large, the matrix  $P$  becomes similarly large, so it becomes expensive to store the matrix and to multiply a vector by  $P$ . Download the script `estatic.m` and open it in your favorite editor. Several lines of this script have been left blank, for you to fill in the details.

- (a) Fill in the matrix  $P$ : the  $(i, j)$  entry should equal the electrostatic potential induced at destination point  $i$  due to a unit charge at source point  $j$ . Equation (1) tells you the energy of the interactions between  $i$  and  $j$ , which equals the charge at  $i$  times the potential at  $i$  due to  $j$ . Check that  $P$  is symmetric.
- (b) Use MATLAB's `eig` function to decompose  $P$  into a matrix  $V$  of eigenvectors and a diagonal matrix  $D$  of the eigenvalues. Check that

$$P = VD V^{-1} \quad (5)$$

and that the columns of  $V$  are orthonormal (that is,  $V_i^T V_j = 1$  if  $i = j$ , or 0 otherwise). Hence  $V^{-1} = V^T$  (in other words,  $V$  is *unitary*). Look at the entries of  $D$  when the separation distance is small and when the separation distance becomes very large. How does the spectrum (the set of eigenvalues) change?

- (c) Imagine that one of the diagonal entries in  $D$  was actually zero. Then no product  $Pq$  would ever have any component along the direction of the corresponding eigenvector. Verify this, if you like; refer to Trefethen and Bau for more details. Similarly, if an eigenvalue  $\lambda_i$  were extremely small, say  $O(\epsilon_{\text{machine}})$  (the computer's precision) then most likely  $Pq$  would have an extremely small component in the direction  $v_i$  (the corresponding eigenvector).

If we know "ahead of time" that we will have a few dominant eigenvalues and everything else will be very small, we might design an approximation to  $P$  in the following way: throw away all the eigenvalue/eigenvector pairs when the eigenvalue is "small enough." Then, to multiply by  $\hat{P}$ , our *low-rank approximation* to  $Pq$  is

$$Pq \approx \hat{P}q = \hat{V} \hat{D} \hat{V}^T q \quad (6)$$

where  $\hat{V}$  denotes the eigenvectors corresponding to the largest  $k$  eigenvalues and  $\hat{D}$  is the  $k$  by  $k$  matrix of those eigenvalues. Calculate  $\hat{P}$  for  $k = \{1, 2, 8\}$ .

- (d) Vary the separation between the clusters and plot the relative errors  $\|Pq - \hat{P}q\|$  as the separation increases.
- (e) Assume there are  $n$  charges in each cluster. How much memory is required to store  $P$ ? How many floating point operations are required to find the product  $Pq$ ?
- (f) If we don't store  $P$  but instead a rank  $k$  approximation  $\hat{P}$ , how much memory is needed? How many floating point operations are needed to calculate  $\hat{P}q$ ? In the limit as  $n$  becomes very large, how do  $P$  and  $\hat{P}$  compare?